

Quasi-Static and Dynamic Large Strain Shear-Tension Testing

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Abstract A new shear specimen is designed, evaluated and tested quasi-statically and dynamically. The specimen consists of a long cylinder having a horizontal gauge section created by two diametrically opposed semi-circular slots machined parallel to the longitudinal axis. This geometry imposes a rather uniform stress state, close to pure shear in the gauge section. Quasi-static and dynamic tension-shear tests up to a strain rate of 10^4 1/s were carried out on 1020 cold-drawn steel specimens. The obtained stress–strain curves and ductility were validated numerically. The new specimen can be used to study the shear mechanical characteristics of a material using tensile testing.

Keywords Large strain · Shear-tension · Quasi static testing · Dynamic testing · Steel 1020

Introduction

The true stress–strain curves and ductility of materials are needed for accurate numerical simulations in diverse areas such as ballistic impact, crashworthiness, machining, forming, frictional effects, erosion and seismic faulting [1, 2]. These properties are dependent on strain-rate, temperature, triaxiality and Lode parameter, not always available in the literature. To generate these properties, extensive research work has been conducted [3–11].

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A wide spectrum of test specimens have been developed to examine and quantify the shear propensity of materials [1, 11–13], but only a few are used with conventional quasi-static and dynamic testing machines alike. Among those, one should mention the hat-specimen [14–17], the shear-compression disc [18, 19], and others for sheet metals [20, 21].

In this study a new shear specimen is designed, evaluated and tested quasi-statically and dynamically. The shear specimen consists of a long cylinder having a horizontal gauge section created by two diametrically opposed semi-circular slots that are machined parallel to the longitudinal axis. This geometry allows for tests to be carried out both with a conventional uniaxial and Tensile Hopkinson Split Bar method (SHTB) [22]. The specimen and its data reduction is similar to the previously proposed shear compression specimen (SCS) [23, 24], and shear tension specimen (STS) [25].

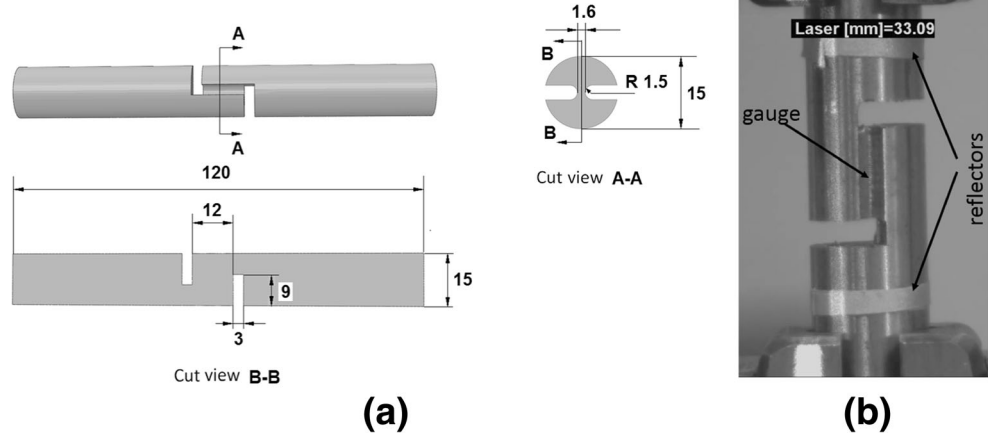
Materials and Methods

Quasi-Static Tensional Tests

Six specimens (Fig. 1(a)) were tested in tension in a servo-hydraulic machine (MTS) with a crosshead velocity of 0.5 mm/min. The extensional displacement (δu) and applied load (F) were measured. The extensional displacement was measured by a laser extensometer and is the distance between two reflectors located on both side of the specimen gauge at an initial gap of ~ 33 mm. The applied load was measured by the MTS load cell. The mounted specimen during test is shown in Fig. 1(b).

The experimental results are summarized in Fig. 2(a). The applied load in Fig. 2(a) is divided by the mid-section area of the gauge. There is a very good repeatability of the results

Fig. 1 (a) The quasi static specimen dimensions. (b) A mounted specimen in the MTS servo hydraulic machine



although specimen number 1 failed earlier, while specimen number 4 underwent slightly higher stresses.

The data reduction technique involves two phases of numerical simulations:

- 1) Pre-test simulation.
- 2) Verification.

Abaqus standard [26] is used for the pre-test simulations which are quasi-static, 3D, and use elastic-plastic material model with Mises plasticity. The applied material properties are taken from literature or previous test results with other specimens. The applied displacement (ΔU) is converted into averaged equivalent plastic strain on the mid-cut section, using $\hat{\epsilon}_p = f_1(\Delta U)$. The applied load (F) is converted to the averaged Mises stress on the mid-cut section using $\hat{\sigma} = g(\hat{\epsilon})F$. The detailed formulas and coefficients are given in appendix 1.

The resulting stress-strain curve is shown in Fig. 2(b) by the dotted (blue) curve. This curve was extended linearly because there are points on the mid cut section that experience higher strains than the average strain which is predicted by the data reduction formulas. This curve serves as a first guess for the validation process.

The same numerical simulations are conducted during the validation phase. Slight changes to the material input of the predicted $\sigma_{eq} - \epsilon_p$ curve of Fig. 2(b) are added until a good agreement between the experimental and the numerical load-displacement curve is obtained. A ductile failure strain criterion [27] is added. This strain, ϵ_p^f , is changed systematically until a good agreement for the location of the failure on the load-displacement (Fig. 2(a)) is obtained. The validated $\sigma - \epsilon_p$ is shown in Fig. 2(b) by the red line. When this $\sigma - \epsilon_p$ relationship is used as input in the numerical simulations, a very good agreement between the numerical $F/(Lt) - \Delta U$ to the experimental ones is achieved. The

numerical $F/(Lt) - \Delta U$ for three different failure strain: $\epsilon_p^f = 1.4, 1.5$ and 1.6 are shown in Fig. 2(a). It can be observed that $\epsilon_p^f = 1.5$ is an appropriate value for the ductility.

The triaxiality (t_r) and the Lode parameter (μ) are defined by equations (1–2) [28–30] as:

$$t_r = -\frac{p}{\sigma_e} \quad (1)$$

$$\mu = \frac{2\sigma_2 - \sigma_1 - \sigma_3}{\sigma_1 - \sigma_3} \quad (2)$$

Where the pressure (p) and the equivalent Mises stress (σ_e) are given in equations (3–4)

$$p = -\sigma_m = -\frac{1}{3}\sigma_{ii} = -\frac{1}{3}(\sigma_1 + \sigma_2 + \sigma_3) \quad (3)$$

$$\begin{aligned} \sigma_e &= \sqrt{\frac{3}{2}s_{ij}s_{ij}} \\ &= \sqrt{\frac{1}{2}[(\sigma_1 - \sigma_2)^2 + (\sigma_2 - \sigma_3)^2 + (\sigma_3 - \sigma_1)^2]} \end{aligned} \quad (4)$$

$$s_{ij} = \sigma_{ij} - \sigma_m I \quad (5)$$

The stresses $\sigma_1 > \sigma_2 > \sigma_3$ are the principal stresses.

Figure 2(c) shows the variation of the *averaged* triaxiality and Lode (t_r, μ) parameters during the plastic deformations. The triaxiality ranges from $0.015 < \hat{t}_r < 0.165$ with an average of 0.083 . The Lode parameter ranges from $-0.3145 < \hat{\mu} < -0.0425$ with an average of -0.155 . The low triaxiality and Lode parameter indicate that the stress situation during plastic deformation is close to pure shear.

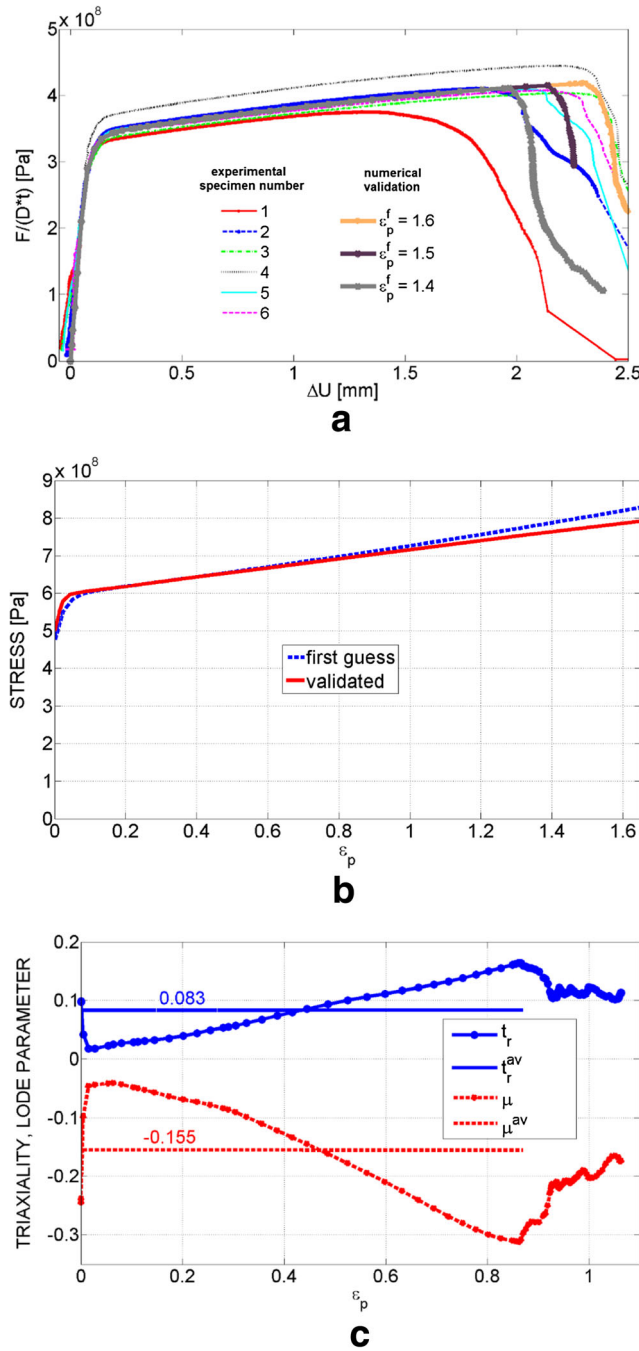


Fig. 2 (a) The six experimental $F/(Lt) - \Delta U$ in comparison to the numerical $F/(Lt) - \Delta U$ which is obtained by using the validated $\sigma_{eq} - \epsilon_p$ for 3 different failure strain: $\epsilon_p^f = 1.4, 1.5$ and 1.6 . (b) The validated $\sigma - \epsilon_p$ in comparison to the predicted $\sigma - \epsilon_p$. (c) The variation of the averaged values on the mid cut section of the triaxiality and Lode parameter during loading. Note the average values during loading

Dynamic Tension Tests

Five shear specimens (Fig. 3(a)) were tested in a split Hopkinson tension bars apparatus (SHTB) [22], made of hardened C300 maraging steel tension bars. A sketch of

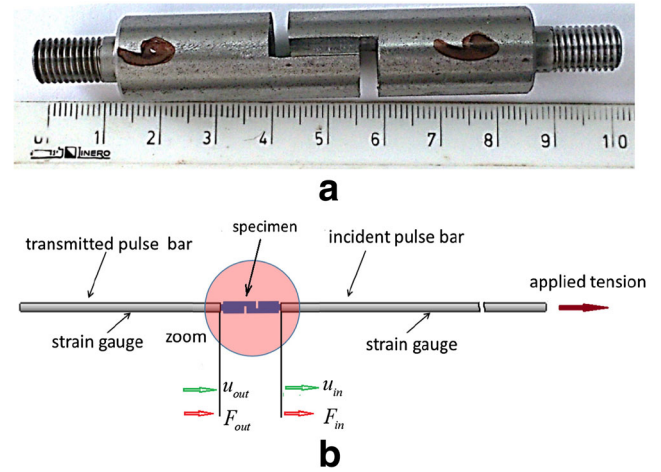


Fig. 3 (a) The shear specimen for dynamic testing. (b) A schematic draw of the SHTB showing the location of the specimen and the location of the obtained forces and displacements

the SHTB is shown in Fig. 3(b). The geometry of the specimens is similar to the quasi-static specimens (Fig. 1) with $D = 15$ mm but shorter length. The specimen is screwed on both sides to the bars using 5/16 UNF threads with 10 mm length.

The experimentally obtained F_{in} and F_{out} versus the applied displacement Δ ($\Delta = u_{in} - u_{out}$) are shown in Fig. 4(a). A high degree of repeatability of the experimental results is noticeable. It can be observed that during the early stage of the pulse, up to approximately $\Delta u = 0.97$ mm, corresponding to ~ 80 μs , there is no dynamic load equilibrium; hence properties obtained in this time (or Δ) regime are not accurate.

The dynamic flow behavior and ductility of the shear specimen is obtained by hybrid numerical-experimental technique similar to the STS. We conducted a 3D non-linear dynamic simulation of the experimental setup using Abaqus explicit [26]. An elastic material model was used for the SHTB bars (Appendix 2). An elastic-plastic Johnson–Cook (JC) material model was used for the 1020 steel (Appendix 2) specimen together with a ductile failure criterion [27].

We use here the JC material model [30] for identifying the dynamic $\sigma - \epsilon_p$ of the material. The flow stress σ_{eq} is assumed to be of the form:

$$\sigma_{eq} = [A + B(\epsilon_p)^n] \left[1 + C \log \left(\frac{\dot{\epsilon}_p}{\dot{\epsilon}_r} \right) \right] [1 - \Theta^m] \quad (6)$$

with the usual definition of this model's parameters (Appendix 2).

We used the quasi-static curves shown in Fig. 2(b) to determine the parameters A, B and n: $A = 500$ MPa, $B = 227$ MPa and $n = 0.41$. The parameter m was set to 1.0

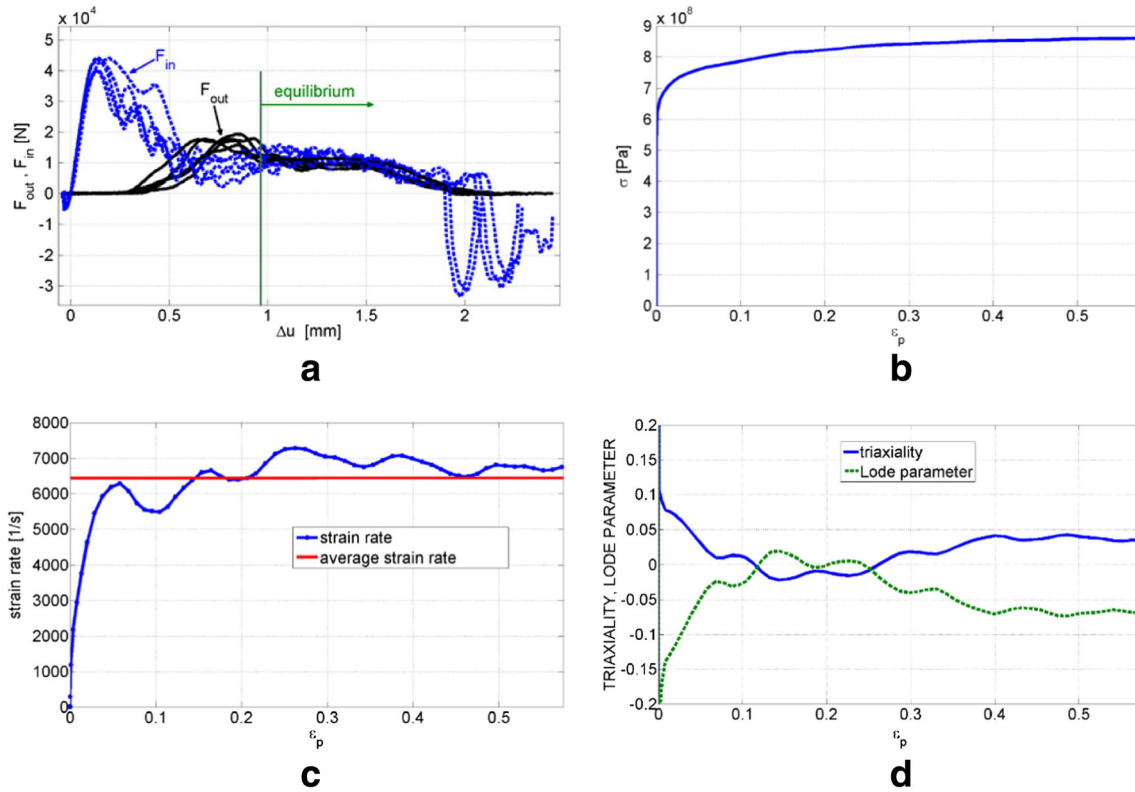


Fig. 4 (a) The experimentally obtained F_{in} and F_{out} versus Δ . (b) The averaged σ - ϵ_p in mid-section of the shear specimen gauge during dynamic shear. (c) The averaged strain rate vs. ϵ_p in mid-section of the shear specimen gauge during dynamic shear showing an average value of ~ 6500 1/s during impact. (d) The averaged triaxiality and Lode parameter vs. ϵ_p in mid-section of the shear specimen gauge during the dynamic shear

– common value for steel [31]. We assumed $\dot{\epsilon}_p^r = 1$ and systematically changed C until a good agreement is obtained between the numerical and experimental F_{in} and F_{out} .

Figure 5(a) shows the converged ($C = 0.04$, $\epsilon_p^f = 0.8$) numerical $F_{out}(\Delta)$ in comparison to five experimental curve. A good agreement is achieved for $\Delta > 1$ mm for which dynamic equilibrium exists (see Fig. 4(a)). The “overshoot” of the experimental $F_{out}(\Delta)$ for $\Delta < 1$ mm is probably due to inertial and thread-related effects [25], keeping in mind that the threads were not modelled explicitly. Numerical tests in which the material was assumed to have higher strength and a drop of strength were carried out. These tests denied this possibility since the resulting $F_{out}(\Delta)$ and $F_{in}(\Delta)$ were completely different from the measured ones. A few values of failure strain were tested numerically and the value of $\epsilon_p^f = 0.8$ cause a best fit between the numerical and experimental F_{out} . The dynamic shear ductility is therefore 0.8. A good agreement between the converged numerical $F_{in}(\Delta)$ to the five experimentally measured ones is shown in Fig. 5(b).

A summary of the JC parameters of cold rolled 1020 steel as well as the elastic properties of the bars is given in Appendix 2.

The averaged σ - ϵ_p on the mid-section of the gauge during the impact is shown in Fig. 4(b). This curve is affected by the evolving strain rate and temperature during the impact. The corresponding strain rate is shown at Fig. 4(c). The strain rate reaches a maximum value of 7400 1/s at $\epsilon_p^f = 0.25$ and has an average value of 6500 1/s. The averaged triaxiality and Lode parameter vs. ϵ_p on the mid-section of the gauge are shown in Fig. 4(d). It can be observed that their absolute value is lower than 0.05 for $\epsilon_p > 0.05$, indicating a dynamic stress situation which is very close to pure shear.

Summary and Conclusions

A new shear specimen, inspired from [1, 21, 24, 25], is *designed, evaluated and tested in tension quasi-statically and dynamically*. The specimen geometry imposes a stress state in the gauge section that is close to pure shear. For 1020 cold drawn steel, the quasi-static ductility at (absolute value) of triaxiality and Lode parameter of ~ 0.1 is 1.5. The dynamic ductility at strain rate of ~ 6500 1/s was found to be 0.8, hence the high shear strain rate elevates the flow stress but lowers the dynamic ductility.

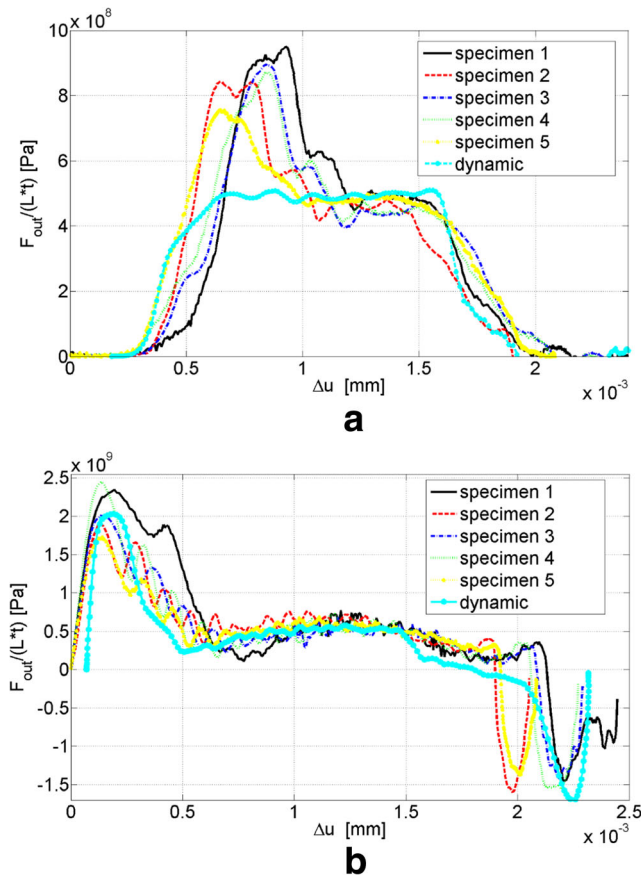


Fig. 5 (a) A comparison between experimental $F_{out}(\Delta)$ and the numerically converged $F_{out}(\Delta)$. (b) A comparison between experimental $F_{in}(\Delta)$ and the numerically converged $F_{in}(\Delta)$

The new specimen can be used to study the shear mechanical characteristics of a material at both low and high strain rates.

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Appendix 1

The data reduction formulas

$$\widehat{\varepsilon}_p = k_1 \frac{\Delta U - \Delta U_y}{L} + k_2 \left(\frac{\Delta U - \Delta U_y}{L} \right)^2 \quad (7)$$

$$\widehat{\sigma} = \left[a e^{b \widehat{\varepsilon}_p} + c e^{d \widehat{\varepsilon}_p} \right] \frac{F}{L \cdot t} \quad (8)$$

In equations (7)–(8) the gauge length is L and the gauge thickness is t . The displacement at which the gauge start to yield is

ΔU_y . The coefficients k_1 , k_2 , a , b , c and d are detailed in Table 1.

Table 1 Coefficients for data reduction

k_1	k_2	a	b	c	d
6.872	-12.0	1.744	-0.04869	-0.05305	-7.706

Appendix 2

Summary of material properties used in numerical analyses.

Table 2 Material properties

	Steel 1020	C300 Managing steel
density $\rho \frac{Kg}{m^3}$	7870	8000
Young's modulus E GPa	205	184
Poisson's ratio ν	0.29	0.3
c m/s	5104	4800
Specific heat $c_p \frac{J}{Kg \cdot C}$	486	—
Inelastic heat fraction β	0.9	—

The Johnson-Cook material model is given by:

$$\sigma_{eq} = [A + B(\varepsilon_p)^n] \left[1 + C \log \left(\frac{\dot{\varepsilon}_p}{\dot{\varepsilon}_p^r} \right) \right] [1 - \Theta^m] \quad (9)$$

Where

$$\Theta = \begin{cases} 0 & \text{for } T < T_r \\ \frac{T - T_r}{T_m - T_r} & \text{for } T_r < T < T_m \\ 1 & \text{for } T > T_m \end{cases} \quad (10)$$

The parameters are summarized in Table 3.

Table 3 JC parameters for cold rolled 1020 steel in dynamic shear at $\sim 10^3$ – 10^4 1/s

A [MPa]	B [MPa]	n	C	m	Tr [C°]	Tm [C°]	$\dot{\varepsilon}_p^r$	ε_p^f ductility
500	227	0.41	0.04	1.0	25	1456	1	0.8

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